

Invasion Percolation and Global Optimization

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Invasion bond percolation (IBP) is mapped exactly into Prim's algorithm for finding the shortest spanning tree of a weighted random graph. Exploring this mapping, which is valid for arbitrary dimensions and lattices, we introduce a new IBP model that belongs to the same universality class as IBP and generates the minimal energy tree spanning the IBP cluster. [S0031-9007(96)00106-8]

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Flow in a porous medium, a problem with important practical applications, has motivated a large number of theoretical and experimental studies [1]. Aiming to understand the complex interplay between the dynamics of flow processes and randomness characterizing the porous medium, a number of models have been introduced that capture different aspects of various experimental situations. One of the most investigated models in this respect is invasion percolation [2], which describes low flow rate drainage experiments or secondary migration of oil during the formation of underground oil reservoirs [1,3].

When a wetting fluid (e.g., water) is injected *slowly* into a porous medium saturated with a nonwetting fluid (e.g., oil), capillary forces, inversely proportional to the local pore diameter, are the major driving forces determining the motion of the fluid. The invasion bond percolation (IBP) model captures the basic features of this invasion process. Consider a two dimensional square lattice and assign random numbers $p_{ij} \in [0, 1]$ to bonds connecting the nearest neighbor vertices x_i and x_j . Here p_{ij} mimic the randomness of the porous medium, corresponding to the random diameter of the pores, and vertices correspond to throats. Invasion bond percolation without trapping is defined by the following steps: (i) Choose a vertex on the lattice. (ii) Find the bond with the smallest p_{ij} connected to the occupied vertex and occupy it. At this point the IBP cluster has two vertices and one bond. (iii) In any subsequent step find the empty bond with the smallest p_{ij} connected to the occupied vertices, and occupy the bond and the vertex connected to it.

The various versions of the model are useful in matching the simulated dynamics to the microscopic effects acting as fluids with different wetting properties and compressibility are considered. Originally introduced to model fluid flow, lately invasion percolation is viewed as a key model in statistical mechanics, investigated for advancing our understanding of irreversible and nonequilibrium growth processes with generic scaling properties [3].

Finding the shortest spanning tree of a weighted random graph is a well known problem in graph theory [4]. Consider a connected nondirected graph G of n vertices and m bonds (links connecting vertices), with costs p_{ij} associated with every bond (x_i, x_j) . A *spanning tree* on

this graph is a connected graph of n vertices and $n - 1$ bonds. Of the many possible spanning trees one wants to find the one for which the sum of the weights p_{ij} is the smallest. A well known example is designing a network that connects n cities with direct city-to-city links (whose length is p_{ij}) and shortest possible total length. This is a problem of major interest in the planning of large scale communication networks and is one of the few problems in graph theory that can be considered completely solved. Since for a fully connected graph with n vertices there are n^{n-2} spanning trees [5], designing an algorithm that finds the shortest one in nonexponential time steps is a formidable global optimization problem.

An efficient algorithm for finding the shortest spanning tree of an arbitrary connected graph G was introduced by Prim [6], and involves the following steps: (i) Choose an arbitrary vertex, x_i . (ii) Of all vertices connected to x_i find the one for which p_{ij} is the smallest, and join x_i and x_j . (iii) At any subsequent step a new vertex is appended to the tree by searching for the bond that has the smallest weight p_{ik} , where x_i belongs to the tree, and x_k does not. Thus bonds that connect already occupied vertices are not eligible for growth. It has been shown by Prim that the tree generated by the previous algorithm is the smallest energy spanning tree for the graph G [7]. Already at this point one can notice the formal similarity between Prim's algorithm and the IBP model discussed above.

In this Letter I show the equivalence of the IBP model with Prim's algorithm for finding the shortest spanning tree of a weighted random graph [6], and explore the consequences of this interesting mapping. For this I introduce an invasion bond percolation model with a local trapping rule (hereafter called IBPO model). At *every time step* the bonds invaded by the IBPO model form the *minimum energy tree* spanning all vertices of the IBP cluster, where energy is defined as the sum of the invaded random bonds. Moreover, the clusters generated by the IBPO model have the same scaling and dynamic properties as the clusters of the standard IBP model. Thus the two models (and Prim's algorithm) belong to the same universality class. Since the IBPO cluster forms a tree (i.e., is loopless), this result implies that loopless IBP belongs to the same universality class as IBP. The cluster

formed by the invaded bonds coincides with the unique solution of the global optimization problem of finding the smallest energy branching self-avoiding walk connecting all vertices of a finite lattice. Furthermore, the IBPO model is computationally more efficient than the IBP model. The above results are *exact* and are valid for *arbitrary dimensions and lattices*.

The difference between the IBP and IBPO models comes in an additional trapping rule [8]: In the IBPO model only bonds connecting vertices of the cluster to empty vertices are eligible for growth (see Fig. 1). Note that in the IBP model there may be bonds eligible for growth, which connect two already occupied vertices (hereafter these are called *trapped* bonds, since an empty bond is trapped between two occupied vertices). In the IBPO model these trapped bonds are not eligible for growth [9].

Consider the invasion process described by the IBPO model, and assume that invasion ends when *all vertices* of a *finite* lattice have been invaded [10]. The energy of the obtained IBPO cluster is defined by $E = \sum p_{ij}$, where the sum goes over all *occupied* bonds.

With these definitions one can prove the following:
 (a) The cluster generated by the IBPO model has the smallest energy of all possible clusters that span all vertices of the lattice. (b) The obtained cluster is independent of the site chosen as the starting point of the invasion process. (c) Defining time as the number of invaded *vertices*, at any time step the vertices invaded by the IBPO model co-

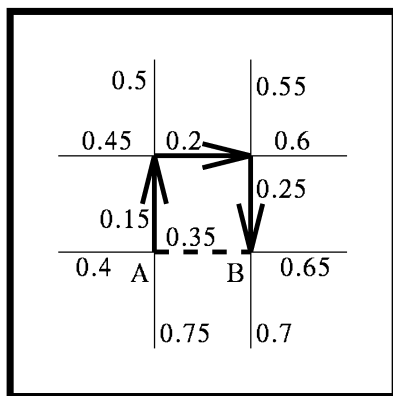


FIG. 1. Definition of the IBP and IBPO models. The figure shows a portion of a two dimensional square lattice with the numbers on the bonds corresponding to the weights p_{ij} . Starting the invasion process from vertex A, both the IBP and IBPO models select the smallest bonds in the order indicated by the arrows. After the third time step the invasion process reaches vertex B. The smallest bond is the one marked with a dotted line, connecting two already occupied vertices, A and B. Such a bond is a *trapped* bond. The IBP model next occupies this trapped bond, without adding any new vertex to the cluster. However, the trapped bond is not eligible for growth in the IBPO model; thus the bond chosen next by the IBPO model is the one with $p = 0.4$. Observe that in the next step the IBP model would choose exactly the same bond. Defining time as the number of occupied *vertices*, the two models occupy the same vertices in exactly the same order.

incide with those invaded by the IBP model, implying that the IBPO and IBP models belong to the same universality class. (d) At any time step the bonds invaded by the IBPO model form the smallest energy tree spanning the vertices of the IBP cluster. (e) The statements (a)–(d) are valid in any dimension and are independent of the lattice.

In the following we discuss (a)–(e) separately.

(a) *Prim's algorithm and IBPO*.—Comparing the definition of the IBPO model and Prim's algorithm, we find that Prim's algorithm is *exactly* the IBPO model acting on the graph G . Since the square lattice, for which the IBPO model is defined, is a particular case of an arbitrary graph, the cluster generated by the IBPO model coincides with the smallest energy spanning tree.

(b) *Uniqueness of the IBPO cluster*.—If the IBPO model selects the smallest energy tree, there is only one such tree, provided that the p_{ij} 's are distinct real numbers, since the chance of having two trees with the same number of bonds and the same energy is zero [11]. Thus starting from any vertex of the lattice one should obtain the same minimum energy cluster.

(c) *Cluster properties*.—A B cluster is the set of *bonds* occupied by the invasion process. Similarly, a V cluster is the set of occupied *vertices*. In percolation and fluid flow one is interested in the scaling properties of the first spanning cluster generated by the invasion algorithm. In particular, it is known that clusters generated by the IBP model are fractal. However, the fractal dimension and in general the scaling exponents may depend on the trapping rule; thus one needs to establish the universality class to which the IBPO model belongs, since it differs from the IBP model in a trapping rule.

Defining time as the number of occupied *vertices*, at any time step the V clusters generated by the IBP and IBPO models are identical [12], the only difference being that within one time step the IBP model may occupy a number of trapped bonds without adding any new vertex to the cluster. The IBPO model with every occupied bond occupies a vertex as well. In conclusion, at any time step the V clusters generated by the two models coincide, provided that we start the invasion process from the same vertex. This implies that the IBP and IBPO belong to the same universality class, and the generated clusters have the same fractal dimension, whose value coincides with the fractal dimension of ordinary percolation [13].

However, not only the static properties, but all dynamic properties measured in terms of the occupied vertices coincide as well. For example, the two models generate exactly the same set of avalanches [14,15] and the growth of the cluster obeys the same dynamic scaling form [15].

(d) *Spanning trees and loopless percolation*.—Next I investigate the relation between the B clusters generated by the two models. The bonds invaded by the IBPO model are a subset of the bonds invaded by the IBP model, i.e., at any time step $N_{IBP}^b \geq N_{IBPO}^b$, where N_{IBPO}^b and N_{IBP}^b are the number of bonds occupied by the IBPO and

IBP models, respectively. According to (a) and Prim's theorem, the bonds invaded by the IBPO model form the smallest energy spanning path connecting the selected vertices. Since the IBP and IBPO models share the same vertices, at every time step the IBPO B cluster is the minimum energy tree spanning all vertices of the IBP clusters. This can be seen in Fig. 2, where the IBP and IBPO clusters are shown simultaneously.

Since the IBPO cluster forms a tree, removing any bond of the IBPO cluster breaks the cluster in two subclusters. This is not true for the IBP model, where by cutting any trapped bond one does not break the cluster (Fig. 2). Since the cluster generated by the IBPO model is a tree, it has no loops. The fact the IBPO and IBP share the same scaling exponents shows that loopless IBP (which is the IBPO model) belongs to the same universality class as IBP, or ordinary percolation. Loopless percolation has been studied in great detail [16], and there is numerical evidence that removing loops does not change the universality class of the percolation model. However, to my knowledge the IBPO model is the first percolation model generating loopless percolation clusters, for which the fact that the loopless model belongs to the same universality class as ordinary (invasion) percolation can be proven exactly.

(e) *Dimension dependence.*—The proof of (a)–(d) does not assume anything specific about the structure of the lattice. Indeed, Prim's theorem applies for an arbitrary weighted graph. Since any regular lattice, in any dimension, is a special case of a random graph, the above results are independent of the nature and dimension of the lattice, proving (e).

Complexity of the IBPO model.—The number of spanning trees on a regular lattice is much smaller than on

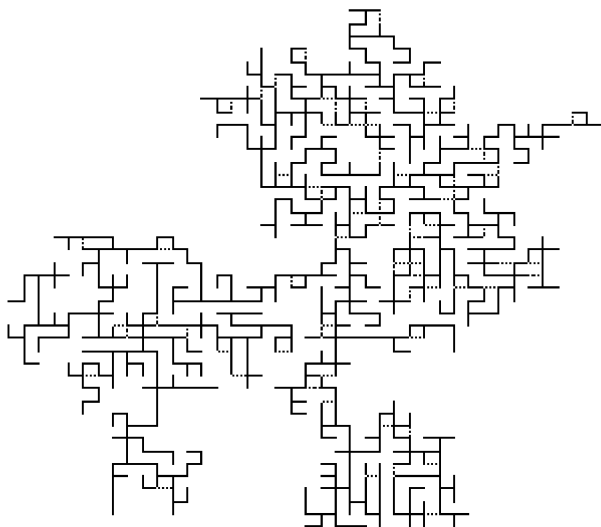


FIG. 2. A particular realization of the IBP and IBPO clusters invading a square lattice simultaneously. The solid bonds form the IBPO cluster. The solid and dotted (trapped) bonds together form the IBP cluster. Note that the IBPO cluster forms a tree (loopless cluster), while every trapped bond leads to a loop to the cluster.

a fully connected graph, but still increases exponentially with n . But the number of computations needed in the simulation of the invasion processes, or the complexity of the IBPO algorithm, is algebraic in n [17]. The most time consuming operation is finding at every time step the bond with the smallest weight eligible for growth. However, since $N_{\text{IBPO}}^b \leq N_{\text{IBP}}^b$, the IBPO model requires equal or less time to run on an arbitrary computer. Figure 3 shows the number of trapped bonds with time $N_{\text{IBP}}^b - N_{\text{IBPO}}^b$. Since the two models belong to the same universality class, using the IBPO model for studying the scaling properties of IBP or ordinary percolation has considerable computational advantages.

In conclusion, I introduced a new bond invasion percolation model that belongs to the same universality class as IBP without trapping, or ordinary percolation. The cluster generated by the IBPO model forms the smallest energy tree spanning the IBP cluster. Exact enumeration, which is the only alternative solution to this global optimization problem, diverges exponentially with the number of vertices in the system. This is the first model, to my knowledge, that through a step-by-step optimization process finds the global minima of the system.

The global optimization problem, to which IBP is shown to be equivalent, connects to another class of problems in statistical mechanics: that of understanding the zero temperature properties of various spatially extended random systems. Since the low temperature behavior is dominated by configurations with the smallest energy, such problems involve finding the minima of certain functions, most often

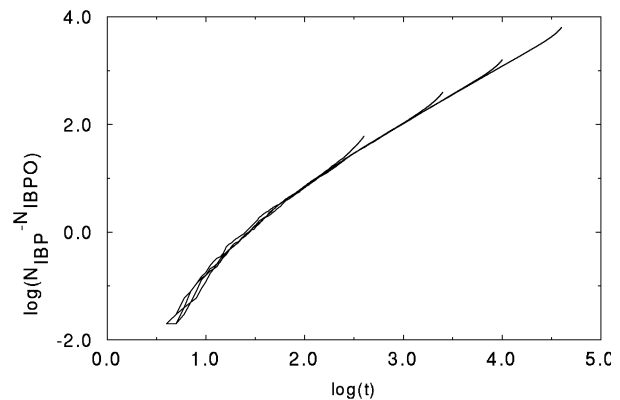


FIG. 3. The difference between the number of bonds occupied by the IBP and IBPO clusters invading a two dimensional square lattice with the same p_{ij} configuration. The horizontal axis corresponds to time, or N_{IBPO}^b , according to the definition of time used in this Letter, while the vertical is $N_{\text{IBP}}^b - N_{\text{IBPO}}^b$. Note that since the difference between the IBP and IBPO models comes in the trapped bonds, $N_{\text{IBP}}^b - N_{\text{IBPO}}^b$ coincides with the number of trapped bonds generated by the IBP model. The four curves correspond to simulations on systems with size $L \times L$, where $L = 20, 50, 100,$ and 200 . An average over 100 runs was taken for each curve. The simulations were stopped after all vertices have been occupied. The asymptotic scaling of the curves suggest a linear behavior.

of a Hamiltonian. Problems in physics that regularly deal with such minimalization procedures range from directed polymers to spin glasses [18], or interface motion in disordered media [19]. The IBPO model provides the minimal energy cluster, implicitly solving a generic problem for a particular random system whose only other solution is exact enumeration.

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